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# HALOGAS: Extraplanar gas in NGC 3198

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## ABSTRACT

We present the analysis of new, deep H I observations of the spiral galaxy NGC 3198 as part of the Westerbork Hydrogen Accretion in Local GALaxieS (HALOGAS) survey, with the main aim of investigating the presence, amount, morphology, and kinematics of extraplanar gas. We present models of the H I observations of NGC 3198. The model that matches best the observed data cube features a thick disk with a scale height of  $\sim 3$  kpc and an H I mass of about 15% of the total H I mass; this thick disk also has a decrease in rotation velocity as a function of height (lag) of  $7\text{--}15 \text{ km s}^{-1} \text{ kpc}^{-1}$  (though with large uncertainties). This extraplanar gas is detected for the first time in NGC 3198. Radially, this gas appears to extend slightly beyond the actively star-forming body of the galaxy (as traced by the H $\alpha$  emission), but it is not more radially extended than the outer, fainter parts of the stellar disk. Compared to previous studies, thanks to the improved sensitivity we trace the rotation curve out to larger radii. We model the rotation curve in the framework of modified Newtonian dynamics (MOND) and confirm that, with the allowed distance range we assumed, fit quality is modest in this galaxy, but the new outer parts are explained in a satisfactory way.

**Key words.** galaxies: halos – galaxies: ISM – galaxies: kinematics and dynamics – galaxies: individual: – galaxies: structure

## 1. Introduction

The last couple of decades have seen a wealth of observational data revealing the presence of material outside the plane of disk galaxies (see Sancisi et al. 2008 and Putman et al. 2012 for a review). This material includes hot X-ray emitting gas (e.g. Tüllmann et al. 2006; Li et al. 2008), ionised gas (Collins & Rand 2001; Rossa & Dettmar 2003), dust (Howk & Savage 1999; Ménard et al. 2010), and neutral hydrogen (Fraternali et al. 2002; Oosterloo et al. 2007; Heald et al. 2011). The material is found to be often (but not always, in the case of the H I) closely associated with star formation in the galaxy disk.

The galactic fountain is an example of a framework that attempts to explain the presence and properties of extraplanar gas: gas is expelled to large distances from the disk by supernova explosions and adiabatic expansion, and then it falls back onto the galaxy disk due to radiative cooling (Shapiro & Field 1976; Bregman 1980). For the conservation of angular momentum, there must be a decrease of the rotation velocity (also known as “lag”) with distance from the plane. However, ballistic models of the interplay between the thin disk and the extraplanar gas (e.g. Fraternali & Binney 2006) show that there must be additional mechanisms governing the kinematics of extraplanar gas, because these simple models obtain lags that are too shallow compared with observations. A possibility is the interaction between the uplifted gas and a corona of slowly rotating hot gas, which would explain the magnitude of the observed

lags (Marinacci et al. 2011). Other explanations for the properties of extraplanar gas involve, for example, thermal instabilities in the corona (Kaufmann et al. 2006; but see Binney et al. 2009), pressure gradients, or the effect of the magnetic field (Benjamin 2002).

Observational evidence for extraplanar lagging H I gas comes from both edge-on and moderately inclined galaxies. Edge-on galaxies allow us to assess the extent of this gas and the magnitude of the lag, whereas in moderately inclined galaxies one can investigate the local connection between star formation and extraplanar gas and study its kinematics in detail. Observationally, in moderately inclined galaxies the very existence of the lag allows the disentangling of extraplanar gas from gas that resides in the plane, e.g. in NGC 2403 (Fraternali et al. 2002), NGC 4559 (Barbieri et al. 2005), and NGC 6946 (Boomsma et al. 2008).

In Heald et al. (2011) we presented the Westerbork Hydrogen Accretion in Local GALaxieS (HALOGAS) survey, a systematic investigation of extraplanar gas in 22 spiral galaxies, using very deep H I observations. The main goal of HALOGAS is to investigate the amount and properties of extraplanar gas. The target galaxies were selected to be nearby and to fall into one of two categories: edge-on galaxies (inclination  $i \geq 80^\circ$ ) and moderately inclined galaxies ( $50^\circ \leq i \leq 75^\circ$ ).

NGC 3198 is one of the moderately inclined galaxies in the HALOGAS sample. It is a late-type spiral that is traditionally considered as one of the benchmark galaxies for rotation curve

**Table 1.** Galaxy NGC 3198 and our observations.

Parameter	Value	Reference
Hubbe type	Sc	1
Adopted distance	$13.8 \pm 1.5$ Mpc	2
<i>B</i> -band magnitude	$10.90 \pm 0.07$	1
Synthesised beam FWHM	$35.2 \times 33.5$ arcsec	3
mJy/beam to Kelvin	0.51	3
On-source time	132 h	3
Velocity resolution	$4.12 \text{ km s}^{-1}$	3
RMS noise	$0.21 \text{ mJy beam}^{-1}$	3

**Notes.** The beam FWHM refers to the cube modelled in this paper.

**References.** 1) HyperLeda; 2) Kelson et al. (1999); Freedman et al. (2001); 3) this work.

studies because its H I kinematics is regular and symmetric, its H I disk is very extended, and its inclination angle (around  $70^\circ$ ) is ideal for deriving the rotation curve. Many studies have focussed partly or completely on the kinematics of NGC 3198, e.g. van Albada et al. (1985), Begeman (1989), Blais-Ouellette et al. (2001), Bottema et al. (2002), de Blok et al. (2008), Sellwood & Sánchez (2010), Gentile et al. (2011). The distance of NGC 3198 has been determined to be 13.8 Mpc using Cepheids (Kelson et al. 1999; Freedman et al. 2001; see also Macri et al. 2001), which are usually considered to be one of the most precise distance indicators.

In the present paper we aim to characterise the presence, amount, and kinematics of any extraplanar gas in NGC 3198, using the most sensitive H I data available and a careful analysis based on a tilted-ring model of the data cube. We also look for a link between the presence of this extraplanar gas and star formation in this galaxy. Additionally, we present detailed kinematic modelling of the observations, as well as a mass decomposition of the rotation curve.

## 2. Data acquisition and reduction

### 2.1. H I data

We summarise here the main points of how the data were acquired and reduced, and we refer the reader to Heald et al. (2011) for a more thorough description of the observations and data reduction. Information about NGC 3198 and the observations presented here is given in Table 1.

The observations of neutral hydrogen (H I) were obtained at the Westerbork Synthesis Radio Telescope (WSRT) in the “maxi-short” configuration for  $10 \times 12$  hours with a bandwidth of 10 MHz divided into 1024 channels and for two linear polarizations. The data were reduced using the software package Miriad (Sault et al. 1995). After flagging, calibration, and continuum subtraction, the data were Fourier inverted to obtain a data cube. When inverting, a Gaussian taper with an image-plane width of 30 arcseconds was applied to the UV data, which were also Hanning smoothed to a velocity resolution of  $4.12 \text{ km s}^{-1}$ . The data cubes obtained in this manner have an rms noise of  $0.21 \text{ mJy beam}^{-1}$ . Avoiding application of the taper gives a higher resolution cube ( $18.9 \times 13.6$  arcsec). However, this cube was not used for modelling because one of the main goals of this paper is to investigate the faint diffuse emission, whose signal-to-noise ratio is higher when applying the taper and thus somewhat degrading the resolution.

The CLEAN deconvolution of the dirty cube was done in two steps. First, a CLEAN mask was defined (for every channel

map) where presence of real emission can be established with confidence. In these regions the CLEAN deconvolution was performed down to approximately  $1\sigma$ . Then, the channel maps were CLEANed on the whole field, down to about  $2\sigma$ . Finally, the channel maps were restored using a Gaussian beam with a FWHM size of  $35.2 \times 33.5$  arcsec.

### 2.2. Optical photometry

We obtained photometric observations of NGC 3198 within the framework of HALOSTARS, which is a deep optical survey of the HALOGAS targets with the Isaac Newton Telescope (INT). NGC 3198 was observed in the  $r'$ -band on February 4, 2010, and the total on-source time amount was 3600 s. A co-added image was obtained using the THELI<sup>1</sup> imaging reduction pipeline (Schirmer et al. 2003; Erben et al. 2005). After the overscan- and bias correction, the images were flat-fielded, super-flat-fielded and defringed. We also masked non-Gaussian noise features such as cosmic rays. Using SCAMP (Bertin 2006), the resulting images were photometrically calibrated, background subtracted and co-added, taking the sky background variation in the individual chips into account and using the 2MASS catalogue to solve for astrometric distortion and for the relative sensitivity of the chips. Using the SDSS (DR8) catalogue, we determined the zero point to be  $24.67 \pm 0.02$  mag. From an rms noise of  $0.023 \text{ ADU/s}$  we thus find a  $1\sigma$  level of 26.4 mag over 1 square arcsec. To enhance faint extended features, the image was convolved with a Gaussian with a FWHM 3 arcsec.

### 2.3. H $\alpha$ imaging

The H $\alpha$  observations of NGC 3198 were made at the Kitt Peak National Observatory (KPNO) as part of an observing run aimed at obtaining wide-field optical images of the HALOGAS sample to investigate the relation between extraplanar gas and star formation. We observed over four nights in January 2012 using the Mosaic 1.1 instrument on the 4-meter telescope. We used an H $\alpha$  filter centred at  $6574.74 \text{ \AA}$  and a FWHM of  $80.62 \text{ \AA}$ , which includes [NII] lines at 6548 and  $6584 \text{ \AA}$ . We exposed in the five-dither pattern with six minutes per exposure for a total of 30 min. We also acquired *R*-band images (ten minutes total) for continuum subtraction.

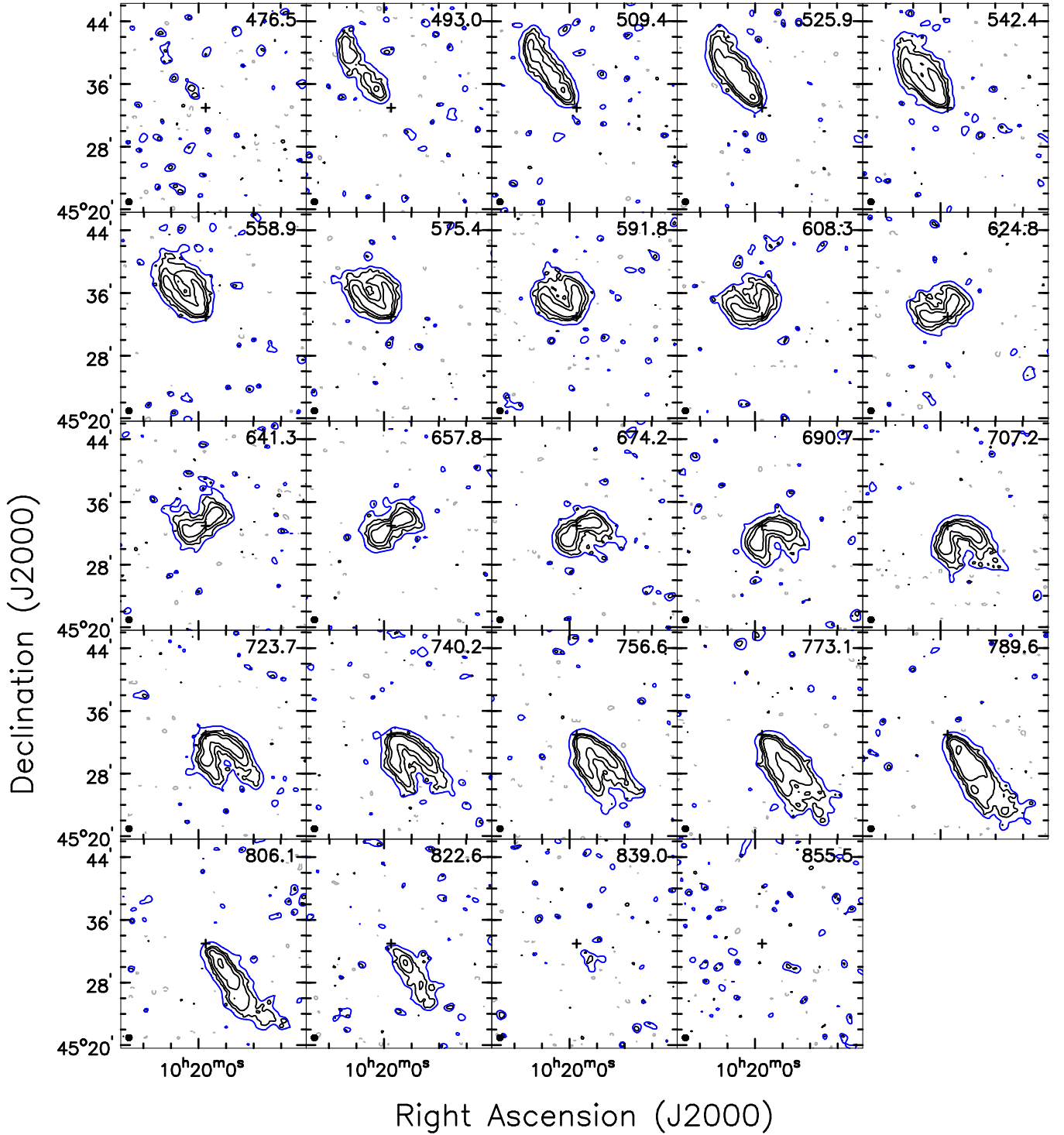
## 3. Observed H I data

The data cube is shown in Fig. 1; for the lowest contour we also plotted the low-resolution (60 arcsec) cube. The channel maps show a regularly rotating disk and some asymmetries (e.g. the south-west side is slightly more extended).

The total H I map was obtained by computing the 0th moment of a masked version of the data cube, which was determined by masking out spurious emission (i.e. emission that is not above  $2\sigma$  in at least three consecutive channels of the 60 arcsec resolution cube). Figure 2 shows the total H I map superimposed with a HALOSTARS image.

From the primary-beam-corrected data cube we find a total H I flux of  $239.9 \text{ Jy km s}^{-1}$ , which translates into a total H I mass of  $1.08 \times 10^{10} M_\odot$ , assuming a distance of 13.8 Mpc. This H I mass is slightly (6%) higher than the value reported in THINGS (The H I Nearby Galaxy Survey, Walter et al. 2008), which might be due partly to amplitude calibration uncertainties and partly to the fact that we detect some extraplanar gas.

<sup>1</sup> <http://www.astro.uni-bonn.de/~theli/>

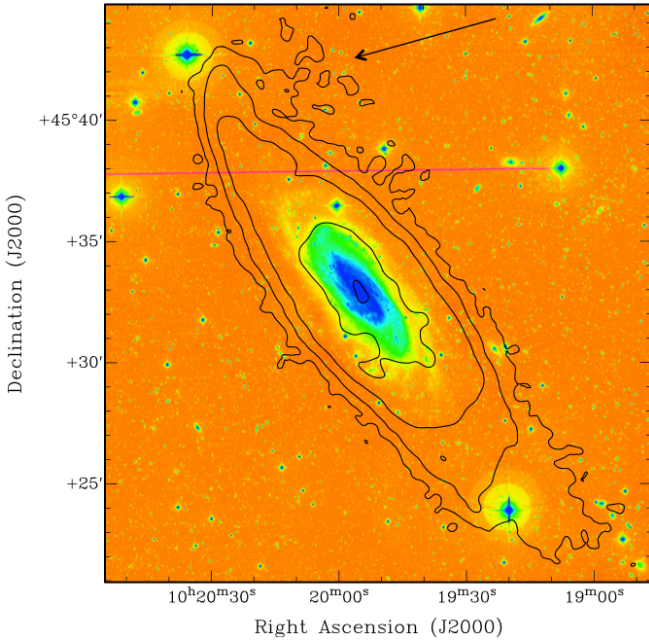


**Fig. 1.** Channels maps of the H I observations of NGC 3198 presented here. The velocity of *each channel map* is indicated in the top right corner of *each panel* (every fourth channel map is shown). The beam is indicated in the bottom left ( $35.2 \times 33.5$  arcsec). The contours are  $-0.42, 0.42$  ( $2\sigma$ ),  $2.1, 10.5, 52.5$  mJy beam $^{-1}$ . The blue contour is the  $0.42$  mJy beam $^{-1}$  ( $2\sigma$ ) level of the  $60$  arcsec resolution cube. Negative contours are dashed grey. The cross indicates the galaxy centre.

In Fig. 2 one can see that our data are more sensitive to faint extended emission than previous observations. At a resolution of  $30$  arcsec, we find meaningful emission down to  $\sim 1 \times 10^{19}$  atoms cm $^{-2}$ . Despite their different angular resolutions, this emission map can be compared with the lowest contour given by Begeman (1989;  $0.5 \times 10^{20}$  atoms cm $^{-2}$ ) and de Blok et al. (2008;  $1 \times 10^{20}$  atoms cm $^{-2}$ ). We note that the latter study focussed

more on having a high spatial resolution than on being sensitive to extended emission. Comparing our map with the one in Begeman (1989), we now confirm the existence of the south-west extension hinted at by Begeman's  $0.5 \times 10^{20}$  atoms cm $^{-2}$  contour (see his Fig. 4). However, the total H I map presented here extends further out, especially in the south-west side of the galaxy, and we tentatively detect some very faint emission to the





**Fig. 2.** Total H I map (contours) superimposed onto a false colour  $r'$ -band HALOSTARS image (see text). Contours are 0.1, 1, 5, and  $15 \times 10^{20}$  atoms  $\text{cm}^{-2}$ . The black arrow indicates the tentatively detected H I described in Sect. 3.

north of the north-eastern side (visible also in Fig. 1, channels at  $493.0 \text{ km s}^{-1}$  to  $558.9 \text{ km s}^{-1}$ ). We estimate the mass of this H I structure (indicated with a black arrow in Fig. 2) to be about  $5 \times 10^6 M_{\odot}$ . We note that this is the only feature that clearly does not belong to the thin or thick disk. In many galaxies there is evidence for more of such features, which indicates some amount of accreting H I.

At a projected distance of about 120 kpc from NGC 3198 we also detect in H I a galaxy pair, identified in the NASA/IPAC Extragalactic Database (NED) as VV 834. One of its two components has a radial velocity from the Sloan Digital Sky Survey of  $566 \pm 256 \text{ km s}^{-1}$ . We find an H I mass of  $2.1 \times 10^7 M_{\odot}$  for the galaxy pair. The centres of the two galaxies are only  $\sim 19$  arcsec apart, and the H I emission is barely resolved, with a total extent of about  $1.5\text{--}2$  arcmin and spreading over  $80 \text{ km s}^{-1}$ , from  $\sim 540 \text{ km s}^{-1}$  to  $\sim 620 \text{ km s}^{-1}$ .

#### 4. Models

To investigate in detail the morphology and kinematics of the H I in this galaxy, we resorted to making models of the whole data cube: it is the most thorough approach that can disentangle subtle differences between various representations of the H I data disk. The uncertainties on the derived parameters are discussed in Sect. 4.1.

The first model we made was based on the parameters derived by de Blok et al. (2008). As can be seen in Fig. 3, the general features of the data cube are reproduced, but these parameters cannot account for the detailed structure of the emission in each channel map. This is partly expected as in the present paper we detect more diffuse emission, which typically has different kinematics.

To improve the model data cube and obtain a better match to the data, we decided to follow the following strategy, based on successive approximations. First, we derived a rotation curve based on a tilted-ring modelling on the intensity-weighted

velocity field (Begeman 1989), leaving the position angle, the inclination angle, the systemic velocity and the rotation velocity as free parameters for each ring. Each side (approaching and receding) was treated separately and the best-fit systemic velocity was found to be  $657 \text{ km s}^{-1}$ . Using the intensity-weighted velocity field can potentially underestimate the inner rotation curve, especially if the velocity gradient is high, as is the case here in the inner regions of NGC 3198. Thus we fixed the position angle and inclination from this fit, and fitted the rotation curve on the velocity field derived using the WAMET method (Gentile et al. 2004), which is a modified version of the envelope-tracing method. For galaxies with a high inclination (such as NGC 3198) or with a poor resolution it is a more reliable choice than the more traditional intensity-weighted mean, which in NGC 3198 might give velocities that are slightly biased towards the systemic velocity. The orientation parameters from the intensity-weighted velocity field were also used to derive the H I surface density profile by averaging the total H I map over ellipses. We used a  $\text{sech}^2$  distribution for the gas in the vertical direction, initially with a scale height of  $0.2 \text{ kpc}$ .

Then, we used these parameters as an initial guess to make a fit of the whole data cube using the Tilted Ring Fitting Code (TiRiFiC) software (Józsa et al. 2007), leaving the inclination, the position angle, and the rotation velocity free for each ring, separately for each side, and the velocity dispersion (including instrumental effects) free as a global parameter. The best-fit velocity dispersion found by TiRiFiC is  $11.7 \text{ km s}^{-1}$ .

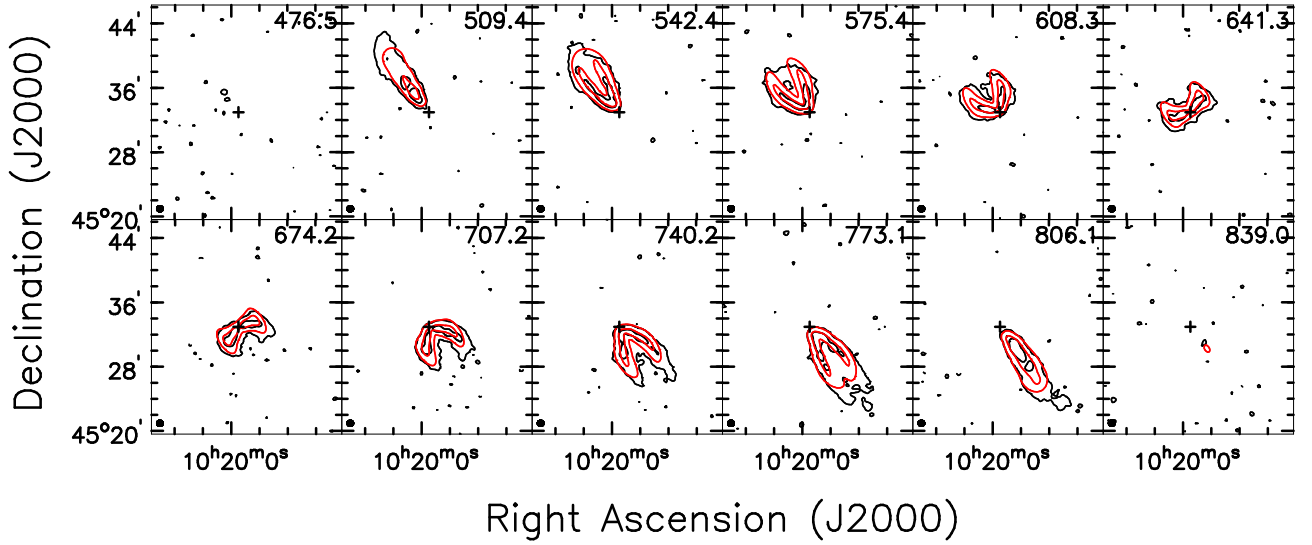
Because of the very high number of degrees of freedom, the parameters of the best-fit model data cube from TiRiFiC tend to have some strong discontinuities and unphysical “jumps” as a function of radius. Also, by construction TiRiFiC uses a  $\chi^2$  minimisation, which naturally (for a constant noise) gives a higher weight to high signal-to-noise regions. However, these regions are not necessarily the most suited to discriminate between models; as a result, the TiRiFiC fitting process has to somehow be manually guided.

Therefore we used this set of parameters as a basis to adjust the parameters manually when necessary, in order to obtain simultaneously a good match of the data cube and a realistic set of parameters. The inclination and the position angle were adjusted based on the total H I map, and the rotation curve was adjusted based on the position-velocity diagram along the major axis. Changes of at most  $5 \text{ km s}^{-1}$  were made to obtain a better match to the observations.

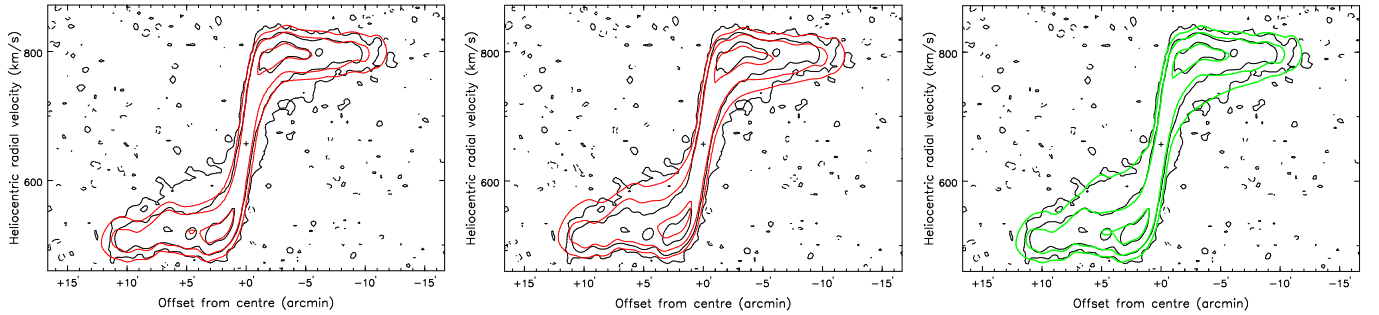
The emission in some channels is skewed (channels  $525.9\text{--}591.8 \text{ km s}^{-1}$  and  $707.2\text{--}740.2 \text{ km s}^{-1}$ , see Fig. 1), without an apparent change in the global position angle and with no sign of position-angle change at other azimuths at these radii.

Spekkens & Sellwood (2007) introduced a formalism to fit velocity fields, focussing on bisymmetric distortions, which they deem more realistic than the more commonly used radial flows. While these two models give similarly good results, Spekkens & Sellwood consider radial flows to be less physically plausible as they have no clear physical origin and are restricted to small non-circular motions due to the epicyclic approximation. Therefore, we interpret the skewness of the emission in channels  $525.9\text{--}591.8 \text{ km s}^{-1}$  and  $707.2\text{--}740.2 \text{ km s}^{-1}$  as a bisymmetric distortion of the velocity field, which is possibly due to an elongation of the potential, ignoring the radial term that is not needed by the present data. In such a model the velocity  $V_{\text{model}}$  at a given position in the galaxy disk is modelled as

$$V_{\text{model}} = V_{\text{sys}} + \sin i [V_{\text{rot}} \cos \theta - V_{2,\text{rot}} \cos(2\theta_b) \cos \theta], \quad (1)$$



**Fig. 3.** Channels maps of NGC 3198: observed (black) and modelled (red), using the rotation curve and inclination parameters from de Blok et al. (2008). Symbols are the same as Fig. 1. For better readability we have plotted fewer contours:  $-0.5$ ,  $0.5$  ( $2.5\sigma$ ), and  $20 \text{ mJy beam}^{-1}$ . Also, compared to Fig. 1, fewer channels are shown.



**Fig. 4.** Position-velocity diagrams of NGC 3198 along the major axis, considering a position angle of  $213^\circ$ : observed (black) and modelled (*left and middle panel*: red; *right panel*: green), derived using the steps outlined in Sect. 4. *Left*: without lagging extraplanar gas. *Middle*: with a single 1 kpc thick disk with a lag. *Right*: with a separate disk of lagging extraplanar gas. Contours are  $-0.3$ ,  $0.3$  ( $\sim 1.5\sigma$ ), 3, and  $30 \text{ mJy beam}^{-1}$ . The cross indicates the galaxy centre.

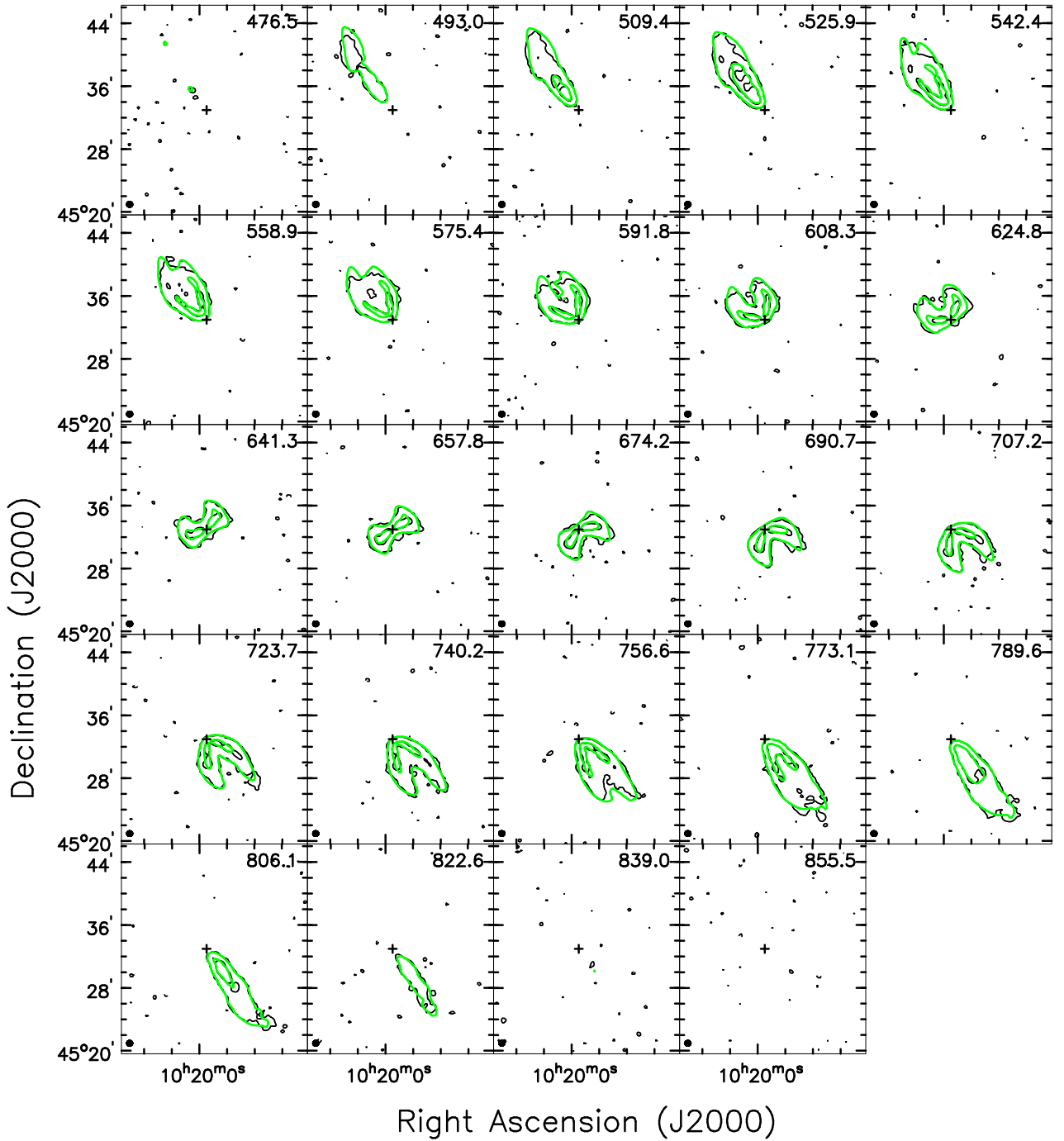
where  $V_{\text{sys}}$  is the systemic velocity,  $V_{\text{rot}}$  is the rotation velocity,  $V_{2,\text{rot}}$  is the  $m = 2$  additional term,  $\theta$  is the azimuthal angle relative to the projected major axis, and if  $\phi_b$  is the angle between the bisymmetric distortion and the projected major axis, then  $\theta_b = \theta - \phi_b$ . Such a distortion, when included in the model, contributes to improving the match between the model and the observations. It turns out that the amplitude  $V_{2,\text{rot}}$  must be  $\sim 15 \text{ km s}^{-1}$  beyond a radius of approximately 510 arcsec and  $\phi_b$  must be  $\sim 45^\circ$  to avoid having skewed emission in the central channels (around the channel at  $657.8 \text{ km s}^{-1}$ ) but to have it at channels  $525.9\text{--}591.8 \text{ km s}^{-1}$  and  $707.2\text{--}740.2 \text{ km s}^{-1}$ . A model with radial motions (limited to a wedge around the major axis to avoid skewed emission in the central channels) gives almost identical results to this bisymmetric distortion.

A position-velocity diagram along the major axis with this model is shown in the left-hand panel of Fig. 4. As clearly visible from the position-velocity diagram, it is necessary to add some lagging gas to reproduce the observations. The first attempt was a single thick disk with a scale-height of 1 kpc and a lag of  $18 \text{ km s}^{-1} \text{ kpc}^{-1}$  in the approaching side and  $8 \text{ km s}^{-1} \text{ kpc}^{-1}$  in the receding side (Fig. 4, middle panel). This model reproduces the observations better than the single thin disk, but it is not satisfactory. By increasing the rotation curve in the inner parts (to match the outermost contours), lagging gas would be missing;

by increasing the amplitude of the lag, the outer regions would not be reproduced correctly.

Therefore we resorted to a two-disk model with a thick disk that rotates more slowly than the main (thin) disk, and with the same orientation parameters and run of surface density. This is shown in the right-hand panel of Fig. 4, where a model that has a lagging thick disk component with a scale-height of 3 kpc is a better representation of the observations. The lag in the model was derived to be  $15 \text{ km s}^{-1} \text{ kpc}^{-1}$  in the approaching side and  $7 \text{ km s}^{-1} \text{ kpc}^{-1}$  in the receding side. The thick disk contains about 15% of the total H I mass and its surface density distribution is a scaled-down version of the total H I density distribution.

The comparison between the final model data cube and the observed one is shown in Fig. 5. Because the lagging extraplanar gas emission is only 1–2% of the peak in each channel map, we investigated whether it could be an artifact due to the CLEAN deconvolution and improper modelling of the dirty beam. To this aim, we built a model data cube based on our final model data cube, only without the thick disk. We then convolved it with the dirty beam, added some Gaussian noise (with an rms chosen to obtain the same rms noise as in the observations), and CLEANed and restored this cube. If the lagging extraplanar gas had been an artifact, in Fig. 6 we would have seen some “pseudo extraplanar gas”. Because we did not, we concluded that the observation of



**Fig. 5.** Channels maps of NGC 3198: observed (black) and modelled (green), using our final model parameters. Symbols are the same as Fig. 1. For better readability we have plotted fewer contours:  $-0.5$ ,  $0.5$  ( $2.5\sigma$ ), and  $20 \text{ mJy beam}^{-1}$ .

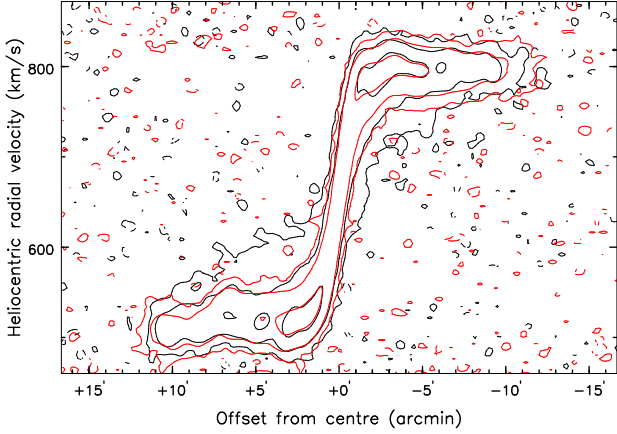
lagging extraplanar gas is genuine and not an artifact due to the dirty beam. This is also confirmed by the hint of a signal from the extraplanar gas that we see in the high-resolution cube (see Fig. 7). The best-fit parameters are plotted in Fig. 8 (only the parameters with a radial dependence): we note that the two sides are quite symmetric.

#### 4.1. Uncertainties on the derived parameters

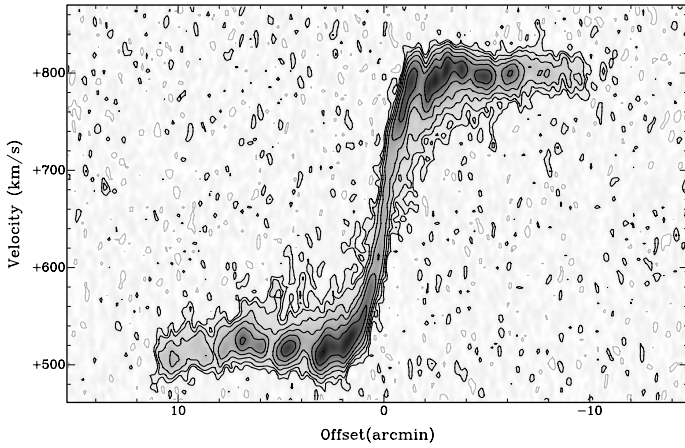
Uncertainties on the parameters that describe the model data cubes were estimated by varying the parameter in question until the model data cube became only marginally consistent with

the observed one, focussing on a particular projection of the data cube or a specific region that is especially sensitive to that parameter.

The inclination and position angle are better derived by comparing total H I maps. A variation in inclination or position angle of only  $2^\circ$  makes the model total H I map barely consistent with the observed one (see Fig. 9 for the inclination). In a similar way, the (global) velocity dispersion is ideally constrained by position-velocity diagrams. It turns out that the high-velocity contours of the position-velocity diagram along the major axis become too spaced out if the velocity dispersion of  $11.7 \text{ km s}^{-1}$  changes by more than  $2 \text{ km s}^{-1}$ .



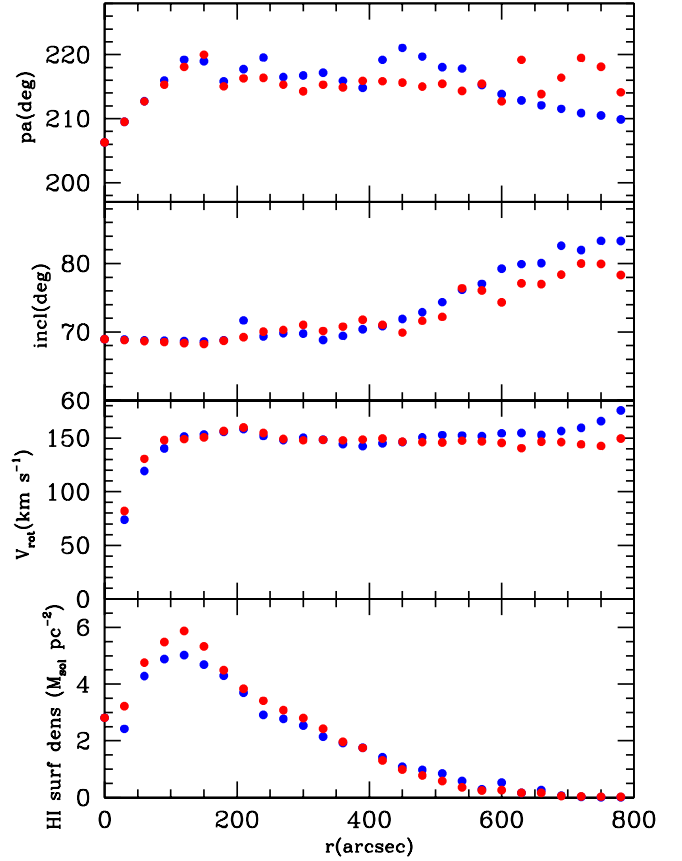
**Fig. 6.** Position-velocity diagram of NGC 3198 along the major axis, considering a position angle of  $213^\circ$ . The observed position-velocity diagram is shown in black and the model is in red. The model shown here is derived with a thin disk only (plus noise), convolved with the dirty beam, CLEANed and restored as described at the end of Sect. 4. Contours are  $-0.3, 0.3$  ( $\sim 1.5\sigma$ ), 3, and  $30 \text{ mJy beam}^{-1}$ .



**Fig. 7.** Position-velocity diagram of the high-resolution cube ( $18.9 \times 13.6 \text{ arcsec}$ ) of NGC 3198 along the major axis, considering a position angle of  $213^\circ$ . Contours are  $-0.29, 0.29$  ( $\sim 1.8\sigma$ ), 0.58, 1.16, ...  $\text{mJy beam}^{-1}$ .

The uncertainties on the rotation curve were assumed to be half the difference between the two sides. To avoid having unrealistically small errors when the velocities of the two sides are close, we also considered a minimum error of  $4.12/\sin(i) \text{ km s}^{-1}$ , where  $i$  is the inclination angle and  $4.12 \text{ km s}^{-1}$  is the channel increment of the data cube we used. The amplitude of bisymmetric distortion can be constrained from the channel maps; we find that it is quite uncertain ( $15 \pm 10 \text{ km s}^{-1}$ ).

How robust are the figures about extraplanar gas? For the thin disk we assumed a scale-height of  $0.2 \text{ kpc}$ , and for the thick disk we tried different values of the scale-height. For a given inclination (constrained from the total H I map, see above), the central channels are particularly sensitive to the scale-height, and we find a value of  $3 \pm 1 \text{ kpc}$ . The relative mass of the extraplanar gas has to be in the range  $10\% - 20\%$ , and the amplitude of the lag is not strongly constrained (NGC 3198 is not edge-on); we find an uncertainty of approximately  $5 \text{ km s}^{-1} \text{ kpc}^{-1}$ . These uncertainties include the degeneracy between scale-height and lag amplitude: a small scale-height can be partly compensated by a large lag value and vice versa. However, a close inspection of the



**Fig. 8.** Radial dependence of the parameters of the best-fit model data cube presented in this paper. Blue represents the approaching (north-east) side and red is the receding (south-west) side. “pa” is the position angle, “incl” is the inclination, “ $V_{\text{rot}}$ ” is the rotation velocity, and “HI surf dens” is the H I surface density.

data cube helps resolve this ambiguity and the above uncertainties take it into account.

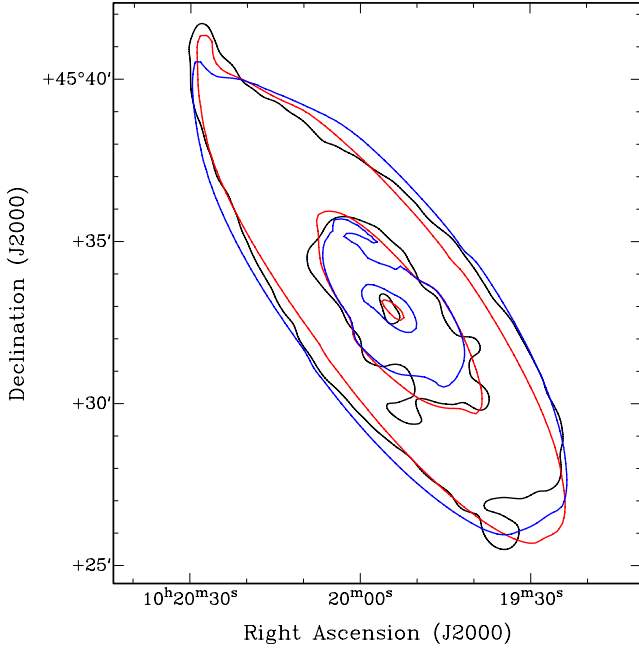
There are, however, two additional uncertainties in the computation of the lag value. First, the lag amplitude might vary with radius (e.g. NGC 891, Oosterloo et al. 2007). We cannot exclude a radially varying lag amplitude with values up to  $10 \text{ km s}^{-1} \text{ kpc}^{-1}$  higher in the inner parts. Second, the value of the lag is also somewhat degenerate with the velocity dispersion of the thick disk; a higher value of the latter can be compensated by a lower value of the lag. However, the velocity dispersion of the thick disk cannot be higher than about  $20 \text{ km s}^{-1}$  (and correspondingly a lag about 25% lower than the above-mentioned values), in which case the emission in the central channels becomes too extended. Lastly, a thick disk with no vertical gradient in the lag but whose rotation velocity is  $20 \text{ km s}^{-1}$  lower than the thin disk yields a model that is almost as good as the model with a vertical gradient in the lag.

## 5. Discussion

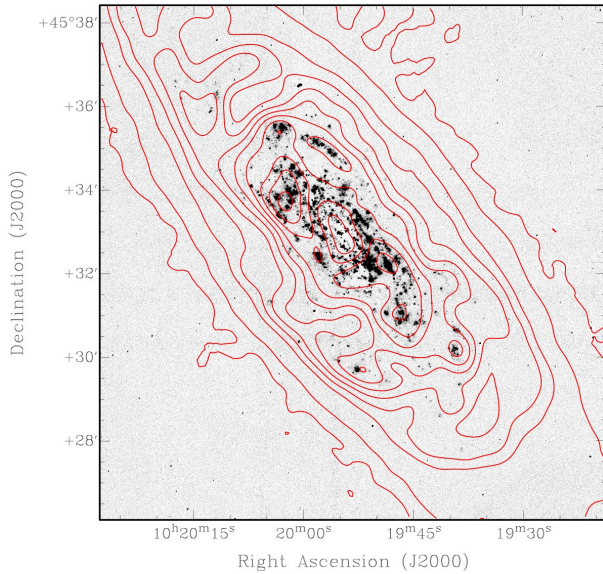
### 5.1. Extraplanar gas

We find that the presence of anomalous gas is required to match the observed and model data cubes. The anomalous gas (defined as the modelled thick disk) has a mass of  $15 \pm 5\%$  of the total H I mass, which is comparable to previous studies that use very deep H I observations to study extraplanar gas, e.g. Zschaechner et al. (2011) and Fraternali et al. (2002). The lag is also a





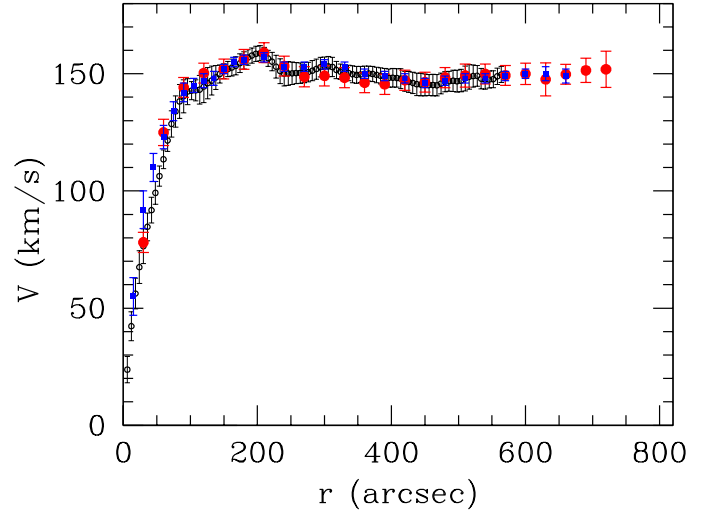
**Fig. 9.** Comparison total H I maps: observed (black), with all the inclination of our best model plus two degrees (red) and minus two degrees (blue). Contours are  $3 \times 10^{20}$  and  $1.5 \times 10^{21}$  atoms  $\text{cm}^{-2}$ .



**Fig. 10.** H $\alpha$  image overlaid with contours of the total H I map. The contour levels are chosen to enhance the central parts of the H I disk:  $1 \times 10^{19}$ , then (1, 5, 7.5, 10, 12.5, 15, 17.5, 20, 22.5, 25)  $\times 10^{20}$ .

required feature of the model, meaning that the extraplanar gas rotates progressively more slowly with increasing vertical distance from the midplane. The value of the lag is comparable with previous estimates, e.g. Oosterloo et al. (2007) and Zschaechner et al. (2011).

The H $\alpha$  image of Sect. 2.3 is shown in Fig. 10. There is a clear correspondence between the brightest H I peaks and H $\alpha$  emitting regions. How does the radial extent of extraplanar gas compare with optical properties? HyperLeda gives an  $R_{25}$  of 194 arcsec, whereas the RC3 gives 255 arcsec. From the position-velocity diagram along the major axis (Fig. 4), there is evidence for extraplanar gas out to 5–6 arcmin (receding side) and at least 6–7 arcmin (approaching side). This can also



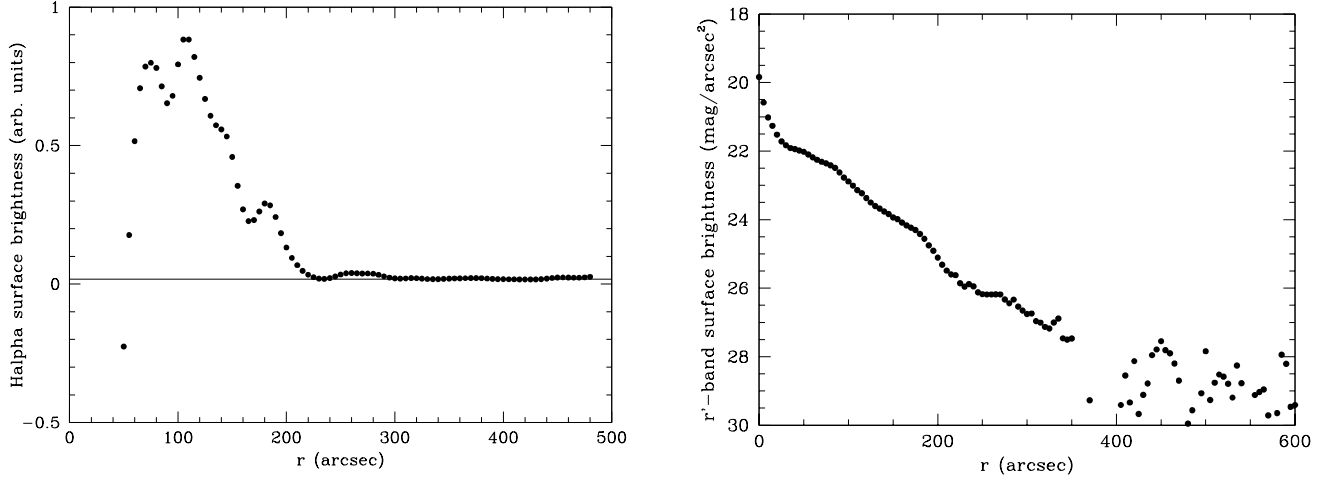
**Fig. 11.** Rotation curve derived here (full red points) compared with the rotation curve derived by de Blok et al. (2008) from the THINGS data (open black points) and by Begeman (1989) using less deep WSRT data (blue squares).

be compared to the extent of H $\alpha$  emission (see Fig. 10 and Fig. 12, left-hand panel) and to our deep  $r'$ -band image from the HALOSTARS survey (Fig. 2, Fig. 12, right-hand panel). Since a precise determination of the extent of the extraplanar gas is problematic because it is quite model dependent, an accurate comparison of the various radial extents (extraplanar gas, stars, H $\alpha$  emission) is hardly possible in a quantitative manner. From the surface brightness profiles (Fig. 12) it appears that the H $\alpha$  emission extends out to  $\sim 3.5$ –5 arcmin (note the central negative values due to slightly imperfect continuum subtraction, which is irrelevant for this paper), whereas the  $r'$ -band image has some very faint emission out to  $\sim 6$  arcmin, comparable to the extent of the extraplanar gas.

Therefore, we conclude that the extraplanar gas extends slightly beyond the actively star-forming body of NGC 3198 but its extent is comparable to the extent of the stellar disk. This extent is also seen in other galaxies with detected extraplanar gas, e.g. NGC 2403 (Fraternali et al. 2002) and NGC 6946 (Boomsma et al. 2008), and it is expected in galactic fountain-type models such as the dynamical models presented by Fraternali & Binney (2006). Regardless of the details of the models, during the initial phases of its motion outside the plane, the gas moves outwards because of the vertical change in the gravitational potential. The origin of extraplanar gas in galaxies is still not unambiguously assessed (as explained more thoroughly in Sect. 1), but its properties in NGC 3198 are similar to those derived for other galaxies, e.g. the Milky Way (Marasco & Fraternali 2011), which were modelled in the framework of galactic fountain-type models.

## 5.2. Rotation curve

The rotation curve we find (Fig. 11) is roughly consistent with the one derived by de Blok et al. (2008) and Begeman (1989). The former focussed more on the high resolution of their data, whereas in the present paper the data have a lower angular resolution but a higher sensitivity to extended emission. In particular, we confirm the velocity decrease by 10–15  $\text{km s}^{-1}$  between 200 and 250 arcsec, and we find a meaningful rotation curve out to (at least) 720 arcsec. The last useful radius is not



**Fig. 12.** Azimuthally averaged surface brightness profiles of the  $H\alpha$  image (left) and our HALOSTARS'  $r'$ -band image (right).

an unambiguous quantity to determine. In order to fully exploit the data, but at the same time to avoid overinterpreting them, we defined (rather conservatively) the last reliable radius as the average between the last radius where the average surface density is above  $1 \times 10^{19}$  atoms  $\text{cm}^{-2}$  (650 arcsec) and the last radius where the tilted-ring fit on the velocity field converged. At these large radii (720 arcsec correspond to 48 kpc for a distance of 13.8 Mpc) there is still no sign of decrease of the rotation curve. We note that in Begeman (1989) the last few points of the rotation curve were derived by assuming the same inclination and position angle as those at a radius of 9.5 arcmin. Moreover, at these radii in the receding side, Begeman's fit to the velocity field is only based on positions outside the major axis.

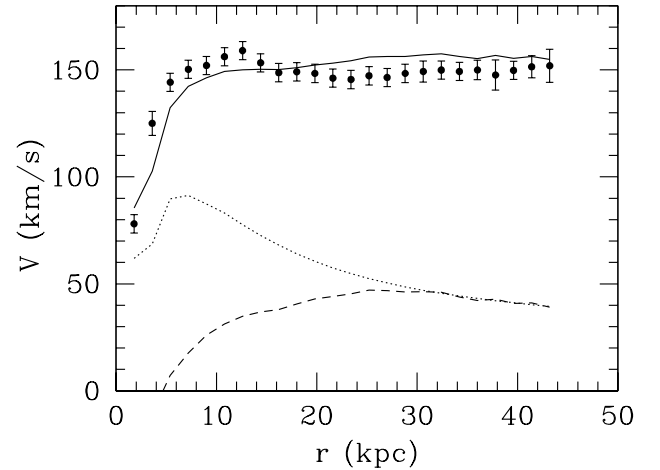
### 5.3. Mass modelling in MOND

The modified Newtonian dynamics (MOND) paradigm was introduced by Milgrom (1983) as an explanation (alternative to dark matter) for the absence of Keplerian decline in the observed kinematics of galaxies. This paradigm has a remarkable predictive power on galactic scales (see e.g. Famaey & McGaugh 2012 and references therein), even though on larger scales it also needs some invisible mass. The galaxy studied in the present paper, NGC 3198, was claimed to show tension with MOND (Bottema et al. 2002; Gentile et al. 2011). It is an ideal case study for MOND because of its inclination (perfectly suited for kinematical studies, see Begeman 1989), its relative closeness, the fact that it is a late-type spiral galaxy with regular and symmetric kinematics, and the accurate determination of its distance using Cepheids (Kelson et al. 1999; Freedman et al. 2001). Indeed, the MOND fit is extremely sensitive to the assumed (or fitted) distance. In MOND, the gravitational acceleration  $g_N$  produced by the visible matter is linked to the true gravitational acceleration  $g$  through the interpolating function  $\mu$ :

$$\mu\left(\frac{g}{a_0}\right)g = g_N, \quad (2)$$

where  $\mu(x) \sim x$  for  $x \ll 1$  and  $\mu(x) \sim 1$  for  $x \gg 1$ . Recently preference has been given to the so-called simple  $\mu$  function:  $\mu(x) = \frac{x}{1+x}$ .

The contribution of the stellar disk was taken from de Blok et al. (2008), where it was derived from  $3.6 \mu\text{m}$  observations (including a correction for a possible mass-to-light ratio gradient, based on the  $J - K$  colour gradient) and the contribution



**Fig. 13.** MOND fit of the rotation curve derived in the present paper, using the distance as a free parameter within the uncertainties of the Cepheids method ( $13.8 \pm 1.5$  Mpc). The dashed line is the Newtonian contribution of the gaseous disks (thin disk and extraplanar gas), the dotted line is the Newtonian contribution of the stellar disk (from de Blok et al. 2008), and the solid line is the best MOND fit.

of the gas from the model data cube discussed above. The thin disk and the thick disk were treated separately (each with its own mass and scale height), and then their velocities were added quadratically.

We analyse here in some more detail the MOND fit with the distance left free within the uncertainty of the Cepheids determination. In comparison with the fit that used the THINGS data (Gentile et al. 2011, Fig. 5), the present MOND fit is of comparable overall quality: see Fig. 13. However, some differences can be seen: in the innermost parts MOND fits the THINGS rotation curve better than the HALOGAS rotation curve, whereas in the outer parts the situation is reversed. We note that MOND fits each rotation curve best where each rotation curve is expected to be the best choice: THINGS in the inner parts (because of the higher resolution) and HALOGAS in the outer parts (because of the better sensitivity to the outer, fainter emission). The best-fit values of the two fits are the same (best-fit  $3.6 \mu\text{m}$  mass-to-light ratio of 0.48 and best-fit distance of 12.3 Mpc, at the lower end of the allowed range). The two fits differ in the shape of the contribution of the gas (apart from, obviously, the rotation curve

itself). With the more extended data from HALOGAS, we account better for the outer surface density profile of the gas and therefore trace better the region where the contribution of the gas start declining. Also, neither of the two fits manages to reproduce the decrease in rotation velocity between 200 and 250 arcsec, which corresponds to the end of the brighter part of the optical disk.

## 6. Conclusions

We have presented new, very deep ( $10 \times 12$  h) H I observations of the spiral galaxy NGC 3198. The observations are part of the Westerbork Hydrogen Accretion in LOcal GALaxies (HALOGAS) survey (see Heald et al. 2011) and go significantly deeper than previous H I observations.

We made careful 3D models of the H I layer in NGC 3198, including not only traditional features such as a variable inclination, position angle, and the rotation curve, but also newer features such as variable rotation speed as a function of distance from the plane, a thicker disk, and a bisymmetric distortion of the kinematics. In this manner we managed to obtain a model that matched the observed data cube very well.

We revealed for the first time in this galaxy the presence of extraplanar gas over a thickness of a few ( $\sim 3$ ) kpc. Its amount is approximately 15% of the total mass, and one of its main properties is that it appears to be rotating more slowly than the gas close to midplane, with a (rather uncertain) rotation velocity gradient in the vertical direction (lag) of  $7\text{--}15 \text{ km s}^{-1} \text{ kpc}^{-1}$ .

Despite the uncertainty on its actual radial extent, the extraplanar gas seems to be slightly more extended than the star-forming part of the galaxy (as revealed by the H $\alpha$  image), but our new deep  $r'$ -band image reveals that there is stellar emission out to the end of the detected extraplanar gas. We detect thin disk H I out to a radial extent that is twice as far as the stellar disk and extraplanar layer. We also detect a faint H I complex beyond the H I disk on the northern side of NGC 3198, with an estimated mass of  $\sim 5 \times 10^6 M_{\odot}$ .

Finally, we make a mass model in the context of MOND of the newly derived rotation curve, which is more extended than previous determinations. The fit quality is modest, similar to previous studies, but the outer parts are explained in a satisfactory way.

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